# A Comparison of the Homotopy Analysis Method and the Homotopy Perturbation Method for the Korteweg-de Vries (K-dV) Equation 

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#### Abstract

In this study, the performance of the Homotopy Analysis Method (HAM) and the Homotopy Perturbation Method (HPM) has been compared for the Korteweg-de Vries (K-dV) equation. The solutions of the K-dV equation by HAM and HPM for three cases have been computed using our MATLAB routine. For a suitable exact solution $l_{2}$-errors has been computed for both the methods. The results show that HPM performs better than HAM for small values of time $t$ and both the results agree well with the exact solution for all the three cases.


Keywords— Homotopy Analysis Method (HAM), Homotopy Perturbation Method (HPM), Korteweg-de Vries (K-dV) equation, MATLAB, partial differential equation (PDE), Auxiliary Parameter, Auxiliary Function, Homotopy Parameter.

## 1 INTRODUCTION

MOST of the natural phenomena are usually expressed by nonlinear partial differential equations (PDEs). The K$d V$ equation, given by

$$
u_{t}+6 u u_{x}+u_{x x x}=0
$$

where $u(x, t)$ represents unknown function, t represents the time, and the subscripts denote partial differentiation. This equation was first used in [21] to represent low-amplitude water wave in shallow, parochial channels such as canals (see [20]). Many researchers have been used different method to solve different types of K-dV equations for various purposes. Comparison of caputo and conformable derivatives for timefractional K-dV equation is studied in [6]. In [8], soliton molecules, nonlocal symmetry and CRE method of the K-dV equation is studied. The K-dV equation is studied for water waves in [10, 11]. The structure of unsteady K-dV model arising in shallow water has been studied in [19]. A comparative analysis of the fractional-order coupled K-dV equations with the Mittag-Leffler Law has been studied in [24]. Fractional forced K-dV equation is studied in [29]. A new localized and periodic solution to a K-dV equation with power law nonlinearity has been studied in [32]. N-soliton solutions and dynamic property analysis of a generalized threecomponent Hirota-Satsuma coupled K-dV equation has been studied in [38]. Moreover, the K-dV equation has now been used to solve variety of problems of different fields including physics, plasma physics and engineering. It is known that most of the nonlinear PDEs do not have analytic solution. For this reason, semi-analytic or numerical solutions are used to solve such problems. There are many semi-analytic methods

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for the solutions of nonlinear PDEs. Among them HAM and HPM are the most popular. HAM was first introduced by Shijun Liao in [34] considering the ideas of homotopy in the general topology. Furthermore, a group of researchers have successfully employed HAM for variety of nonlinearproblems such as: Generalized Sylvester matrix equation with applications is studied in [2]. In [4] two dimensional linear Volterra fuzzy integral equations have been studied. Nonsimilar solution of Eyring-Powell fluid flow and heat transfer with convective boundary condition is studied in [5]. DaveyStewartson equations is studied in [12]. Fuzzy impulsive fractional differential equations have been studied in [25]. In [28] determining the thermal response of convective-radiative porous fins with temperature-dependent properties is studied. SIR epidemic model with Crowley-Martin type functional eesponse And Holling type- II treatment rate has been studied in [30]. Beyond Perturbation: Introduction to the Homotopy Analysis Method is studied in [33]. In [35] $\nabla^{2} u=b(x, y)$ type equations have been studied. And many other types of nonlinear problems are studied.

The HPM is another popular semi-analytic technique for the solution of nonlinear PDEs. It was introduced by Ji-Huan He in [18] using the general concepts of homotopy in topology. HPM has been utilized by many researchers to solve various types of linear and non-linear problems, such as: Mathematical study of diabetes and its complication is studied in [1]. SIR Mumps_Model has been studied in [3]. In [7] nonlinear Volttra partial integro-differential equations have been studied. Nonlinear Schrodinger equations are studied in [ $9,13,14,16]$. Nonlinear Burger equations are studied in [13, 14]. Nonlinear equations arising in heat transforms is studied in [13]. Fractional differential equations: part 1 Mohand transform is studied in [26]. In [27] nonlinear Oscillators is studied. Biological population model has been studied in [31]. In [37] inverse analysis of Jeffery-Hamel flow problem is studied. And many other types of nonlinear problems have been studied.

In this study, we have intended to compare the performance of both HAM and HPM for the K-dV equation and compare both the results with a suitable exact solution for different number of terms and different values of parameters. We have considered three different cases. The results show that HPM performs better than HAM for small values of time $t$, and both the results agreed well with the exact solution for all the three cases.

## 2 MATHEMATICAL BACKGROUND:

2.1 Korteweg-de Vries (K-dV) Equation: In this study we have considered the following K-dV equation,

$$
\begin{equation*}
u_{\mathrm{t}}+6 u u_{\mathrm{x}}+u_{\mathrm{xxx}}=0 \tag{a}
\end{equation*}
$$

with the initial approximation $u(x, 0)=N(N+1) \operatorname{sech}^{2}(x)$, $N>0$, where $u(x, t)$ represents unknown function, t represents the time, and the subscripts in equation (a) denote partial differentiation. Considering $N=1, \quad u(x, 0)=$ $N(N+1) \operatorname{sech}^{2}(x)$ becomes $u(x, 0)=2 \operatorname{sech}^{2}(x)$.

### 2.2 The Homotopy Analysis Method (HAM):

To demonstrate the fundamental concepts of the HAM, we assume

$$
\begin{equation*}
\mathcal{A}[v(t)]-f(t)=0, \quad t \in \Omega \tag{1}
\end{equation*}
$$

be the usual differential equation with boundary condition $\mathcal{B}\left(v, \frac{\partial v}{\partial m}\right)=0, t \in \Gamma$, where $\mathcal{A}$ denotes the nonlinear differential operator, $f(t)$ denotes the known function, $\mathcal{B}$ denotes the boundary operator, $v(t)$ represents the unknown function, $t$ is the time, and $\Gamma$ denotes boundary of the region $\Omega$. The nonlinear differential operator $\mathcal{A}$ can be separated into two parts which are $\mathcal{A}(v)=\mathcal{L}(v)+\mathcal{N}(v)$, where $\mathcal{L}$ represents linear operator, and $\mathcal{N}$ represents non-linear operator. Therefore the equation (1) can be written as follows:

$$
\begin{equation*}
\mathcal{L}(v)+\mathcal{N}(v)-f(t)=0 \tag{2}
\end{equation*}
$$

Using homotopy technique, for a function

$$
\begin{gather*}
\psi: \Omega \times[0,1] \rightarrow \mathbb{R}, \text { we define a homotopy } \\
\widetilde{H}(\psi, p): \mathbb{R} \times[0,1] \rightarrow \mathbb{R} \text { by } \\
\widetilde{H}(\psi(t ; p), p)=(1-p)\left[\mathcal{L}(\psi(t ; p))-\mathcal{L}\left(v_{0}(t)\right)\right]+ \\
p[\mathcal{L}(\psi(t ; p))+\mathcal{N}(\psi(t ; p))-f(t)], \quad \ldots \tag{3}
\end{gather*}
$$

where $p \in[0,1]$ is a homotopy parameter, and $\psi$ is a function of $t$ and $p$.

In [33], Dr. Shijun Liao defined a new type of homotopy
$H(\psi, p): \mathbb{R} \times[0,1] \rightarrow \mathbb{R}$ by introducing a auxiliary parameter $\hbar$ and a auxiliary function $\mathcal{H}(t)$ so that $\hbar \neq 0$ and $\mathcal{H}(t) \neq 0$ as

$$
\begin{array}{r}
H(\psi, p, \hbar, \mathcal{H})=(1-p)\left[\mathcal{L}(\psi(\mathrm{t} ; \mathrm{p}), p, \hbar, \mathcal{H}(t))-\mathcal{L}\left(v_{0}(t)\right)\right]-p \hbar \\
\mathcal{H}(t)[\mathcal{L}(\psi(t ; p), p, \hbar, \mathcal{H}(t))+\mathcal{N}(\psi(t ; p), p, \hbar, \mathcal{H}(t))-f(t)] \tag{4}
\end{array}
$$

Clearly (4) is more general than (3), since (3) is the special case of (4) for $\hbar=-1 \& \mathcal{H}(t)=1$. i.e., $\widetilde{H}(\psi(t, p), p)=$ $H(\psi(t, p), p,-1,1)$.

Then the deformation equation of order zero constructed by Liao [33] is given by

$$
\begin{align*}
H[\psi(t ; p) ; p, \hbar, \mathcal{H}] & =(1-p) \mathcal{L}\left[\psi(t ; p)-v_{0}(t)\right]- \\
p & \hbar \mathcal{H}(t) \mathcal{N}[\psi(t ; p)], \quad \ldots \ldots \ldots \tag{5}
\end{align*}
$$

where $p$ denotes the homotopy parameter with $p \in[0,1], \mathcal{L}$ denotes the auxiliary linear operator with the property that

$$
\begin{equation*}
\mathcal{L}(0)=0 \tag{6}
\end{equation*}
$$

$v_{0}(t)$ denotes the initial approximate solution, and $\mathcal{N}$ denotes the nonlinear operator which is defined as:

$$
\mathcal{N}[\psi(t ; p)]=\mathcal{L}[\psi(t ; p)]+\mathcal{N}[\psi(t ; p)]-f(t)
$$

Substituting $H[\psi(t ; p) ; p, \hbar, \mathcal{H}]=0$ in the equation (5) the zeroth order deformation equation can be written as:

$$
\begin{equation*}
(1-p) \mathcal{L}\left[\psi(t ; p)-v_{0}(t)\right]=p \hbar \mathcal{H}(t) \mathcal{N}[\psi(t ; p)] \ldots \ldots \tag{7}
\end{equation*}
$$

If $p=0$, then we have

$$
\left.\begin{array}{c}
\qquad\left.H[\psi(t ; p) ; p, \hbar, \mathcal{H}]\right|_{p=0}=\mathcal{L}\left[\psi(t ; 0)-v_{0}(t)\right]=0 \\
\text { i.e., } \mathcal{L}[\psi(t ; 0)]=v_{0}(t) \tag{8}
\end{array} \ldots \ldots \ldots\right)
$$

and if $p=1$, then we have

$$
\begin{align*}
& \left.H[\psi(t ; p) ; p, \hbar, \mathcal{H}]\right|_{p=1}=\hbar \mathcal{H}(t) \mathcal{N}[\psi(t ; 1)]=0 \\
& \text { i.e., } \hbar \mathcal{H}(t) \mathcal{N}[\psi(t ; 1)]=0 \quad \ldots \ldots \tag{9}
\end{align*}
$$

Now, it is clear from the equations (6), (8), and (9) that

$$
\psi(t ; 0)=v_{0}(t) \text { and } \psi(t ; 1)=v(t)
$$

Therefore if the homotopy parameter $p$ increases from 0 to 1 , the solution $\psi(t ; p)$ continuously changes from $v_{0}(t)$ to the solution $v(t)$ of the given equation (1). In topology, this type of continuous transformation is known as deformation.
Now, differentiating equation (7) $m$ times w. r. to the homotopy parameter $p$, and putting $p=0$, and finally multiplying them by $\frac{1}{m!}$, we get the deformation equation of order $m$ as follows:

$$
\begin{equation*}
\mathcal{L}\left[v_{m}(t)-X_{m} v_{m-1}(t)\right]=\hbar \mathcal{H}(t) D_{m-1}\left(\bar{v}_{m}\right) \tag{10}
\end{equation*}
$$

where $\bar{v}_{\mathrm{m}}=\left\{v_{0}(t), v_{1}(t), \ldots, v_{m}(t)\right\}$,

$$
\begin{aligned}
& D_{m-1}\left(\overline{\mathrm{v}}_{m}\right)=\frac{1}{(m-1)!} \frac{\partial^{m-1} \mathcal{N}[\psi(t ; p)]}{\partial \mathrm{p}^{m-1}} \text { and } \\
& X_{m}= \begin{cases}0, & \text { when } m \leq 1 \\
1, & \text { when } m>1\end{cases}
\end{aligned}
$$

Since $\psi(t ; p)$ depends on $p \in[0,1]$, by Taylor's theorem we have the series expansion of $\psi(t ; p) \mathrm{w}$. r. to $p$ as

$$
\begin{equation*}
\psi(t ; p)=v_{0}(t)+\sum_{\mathrm{m}=1}^{+\infty} v_{m}(t) p^{m} \tag{11}
\end{equation*}
$$

where $v_{m}(t)=\left.\frac{1}{m!} \frac{\partial^{m} \psi(t ; p)}{\partial p^{m}}\right|_{\mathrm{p}=0}$.
Solving equation (10) we can find $v_{m}(t)$. If the initial approximation $v_{0}(t)$, the auxiliary linear operator $\mathcal{L}$, the auxiliary function $\mathcal{H}(t)$ and the auxiliary parameter $\hbar$ can be chosen properly, then the above series (11) must be convergent at $p=1$.
Then at $p=1$ the series (11) becomes

$$
\psi(t ; 1)=v_{0}(t)+\sum_{\mathrm{m}=1}^{+\infty} v_{m}(t)
$$

Therefore we have

$$
\begin{equation*}
v(t)=v_{0}(t)+\sum_{\mathrm{m}=1}^{+\infty} v_{m}(t) \tag{12}
\end{equation*}
$$

In [33] Liao proved that the series (12) is one of the results of the given equation (1). It should be very significant to assure that at $p=1$, the series (11) must be convergent, on the other hand there is no meaning of the series (12).

### 2.3 The Homotopy Perturbation Method (HPM):

To demonstrate the fundamental concepts of the HPM, we consider a differential equation, which is given by

$$
\begin{equation*}
\mathcal{A}(u)=f(r), r \in \Omega \tag{13}
\end{equation*}
$$

together with

$$
\begin{equation*}
\mathcal{B}\left(u, \frac{\partial u}{\partial m}\right)=0, r \in \Gamma \tag{14}
\end{equation*}
$$

where $\mathcal{A}$ denotes the usual differential operator, $\mathcal{B}$ denotes the usual boundary operator, $f(r)$ denotes a known analytic function, the domain denoted by $\Omega$, and $\Gamma$ denotes the boundary of domain $\Omega$.

Similar to HAM the usual differential operator $\mathcal{A}$ can be separated into two parts as $\mathcal{A}(u)=\mathcal{L}(u)+\mathcal{N}(u)$, where $\mathcal{L}$ stands for the linear operator, and $\mathcal{N}$ stands for non-linear operator in the given differential equation. Therefore equation (13) can be written as

$$
\begin{equation*}
\mathcal{L}(u)+\mathcal{N}(u)-f(r)=0 \tag{15}
\end{equation*}
$$

Using homotopy technique, we can define a homotopy as $w(r, p): \Omega \times[0,1] \rightarrow \mathbb{R}$ and $H(w, p): \mathbb{R} \times[0,1] \rightarrow \mathbb{R}$ satisfying the homotopy equation:

$$
\begin{align*}
H(w, p) & =(1-p)\left[\mathcal{L}(w)-\mathcal{L}\left(u_{0}\right)\right]+p[\mathcal{L}(w)+\mathcal{N}(w)-f(r)] \\
& =0, \quad p \in[0,1], r \in \Omega \tag{16}
\end{align*}
$$

i.e.,

$$
\begin{gather*}
\left.\left.H(w, p)=\underset{\sim}{\mathcal{L}(w)-\mathcal{L}\left(u_{0}\right)+p \mathcal{L}\left(u_{0}\right)+p[\mathcal{N}(w)-f(r)]=0,} \begin{array}{c}
\ldots \ldots \ldots(17
\end{array}\right) . \quad \ldots, 1\right], r \in \Omega,
\end{gather*}
$$

where $p \in[0,1]$ denotes a homotopy parameter and $u_{0}$ denotes an initial approximate solution of the given differential equation (13) satisfying the given boundary conditions.
From equation (16) and (17), we have,

$$
\begin{aligned}
& H(w, 0)=\mathcal{L}(w)-\mathcal{L}\left(u_{0}\right)=0 \\
& H(w, 1)=\mathcal{L}(w)+\mathcal{N}(w)-f(r)=0
\end{aligned}
$$

The changing procedure of $p$ from 0 (zero) to 1 (unity) is only that $w(r, p)$ shifting from $u_{0}(r)$ into $u(r)$, this is said to be homotopy, in topology. Therefore, we have,

$$
\mathcal{L}(w)-\mathcal{L}\left(u_{0}\right) \cong \mathcal{L}(w)+\mathcal{N}(w)-f(r), r \in \Omega
$$

and $\quad w_{0}(r) \cong w(r), \quad r \in \Omega$

In topology, (19) is called deformation. Since $w(r, p)$ depends on homotopy parameter $p \in[0,1]$, by Taylor's theorem we have the series expansion of $w(r ; p) \mathrm{w}$. r. to $p$ as follows:

$$
\begin{equation*}
w=p^{0} w_{0}+p^{1} w_{1}+p^{2} w_{2}+p^{3} w_{3}+\cdots \tag{20}
\end{equation*}
$$

i.e.,

$$
w(r, p)=\sum_{i=0}^{+\infty} p^{i} w_{i}(r)
$$

Consider that this series expansion (20) gives the solutions of the equations (16) and (17).
Setting $p=1$ in the equation (20), we get

$$
w(r, 1)=w_{0}+w_{1}+w_{2}+w_{3}+\cdots
$$

Therefore

$$
\begin{equation*}
u(r)=\lim _{p \rightarrow 1} w(r, p)=w_{0}+w_{1}+w_{2}+w_{3}+\cdots \tag{21}
\end{equation*}
$$

i.e.,

$$
u(r)=\lim _{p \rightarrow 1} w(r, p)=\sum_{\mathrm{i}=0}^{+\infty} w_{i}(r)
$$

which is the solution of the given equation (13).
The perturbation method coupling with the homotopy technique is known as the homotopy perturbation method. This removes the constraint of the usual perturbation method. However, HPM has the full amenities of the usual perturbation method. For most cases, the series (21) is convergent.
However, to trace the rate of convergence on the non-linear operator, the following suggestions have been made by Dr. JH. He [18]:
a) The $\frac{\partial^{2} \mathcal{N}(w)}{\partial w^{2}}$ must be minimal so the parameter could be quite large, i.e., $p \rightarrow 1$.
b) The norm $\left\|\mathcal{L}^{-1} \frac{\partial \mathcal{N}}{\partial w}\right\|<1$ so that the series converges.

## 3 NUMERICAL SCHEME OF THE K-dV EQUATION:

3.1 Numerical Scheme by the HAM: Consider the K-dV
equation (a) with the conferred initial approximation

$$
u(x, 0)=2 \operatorname{sech}^{2}(x)
$$

Then the $\mathrm{m}^{\text {th }}$ order deformation equation is

$$
\begin{equation*}
\mathcal{L}\left[u_{m}(x, t)-X_{m} u_{m-1}(x, t)\right]=\hbar \mathcal{H}(x, t) D_{m-1}[\mathcal{N}(u(x, t))], \tag{22}
\end{equation*}
$$

where $\mathcal{L}(u)=\frac{\partial u}{\partial t}, \mathcal{N}(u)=\frac{\partial u}{\partial t}+6 u \frac{\partial u}{\partial x}+\frac{\partial^{3} u}{\partial x^{3}}$,

$$
x_{m}=\left\{\begin{array}{l}
0, \text { when } \mathrm{m} \leq 1 \\
1, \text { when } \mathrm{m}>1
\end{array} \text { and } u_{0}=u(x, 0)=2 \operatorname{sech}^{2}(x)\right.
$$

Then we have,

$$
\left(u_{m}\right)_{t}-X_{m}\left(u_{m-1}\right)_{t}=\hbar \mathcal{H} D_{m-1}[\mathcal{N}(u(x, t))]
$$

Integrating both sides from 0 to 1 with respect to $t$, we get,

$$
\begin{align*}
& u_{m}-\mathcal{X}_{m} u_{m-1}=\int_{0}^{\mathrm{t}} \hbar \mathcal{H} D_{m-1}[\mathcal{N}(u(x, t))] \mathrm{dt} \\
& \Rightarrow u_{m}=\mathcal{X}_{m} u_{m-1}+\int_{0}^{\mathrm{t}} \hbar \mathcal{H} D_{m-1}[\mathcal{N}(u(x, t))] \mathrm{dt} \tag{23}
\end{align*}
$$

Now, $D_{m-1}[\mathcal{N}(u(x, t))]=D_{m-1}\left[u_{t}+6 u u_{\mathrm{x}}+u_{\mathrm{xxx}}\right]$

$$
\begin{aligned}
& =D_{m-1}\left[u_{t}\right]+D_{m-1}\left[6 u u_{\mathrm{x}}\right]+D_{m-1}\left[u_{\mathrm{xxx}}\right] \\
& =\left(u_{m-1}\right)_{t}+6 \sum_{i=0}^{m-1} u_{i} \cdot\left(u_{m-1-i}\right)_{x}+\left(u_{m-1}\right)_{\mathrm{xxx}}
\end{aligned}
$$

where $D_{m-1}\left[6 u u_{\mathrm{x}}\right]=6 \sum_{\mathrm{i}=0}^{\mathrm{m}-1} \mathrm{u}_{\mathrm{i}} \cdot\left(\mathrm{u}_{\mathrm{m}-1-\mathrm{i}}\right)_{\mathrm{x}}$.

Then the equation (23) can be written as

$$
\begin{align*}
u_{m}=\mathcal{X}_{m} u_{m-1} & +\hbar \mathcal{H} \int_{0}^{\mathrm{t}}\left[\left(u_{m-1}\right)_{t}+6 \sum_{i=0}^{m-1} u_{i} \cdot\left(u_{m-1-i}\right)_{x}\right. \\
& \left.+\left(u_{m-1}\right)_{x x x}\right] d t \tag{24}
\end{align*}
$$

Putting $m=1,2,3, \cdots$, respectively in the above equation (24), we can find $u_{1}, u_{2}, u_{3}, \cdots$, as follows:

$$
\begin{aligned}
& u_{1}=-2^{4} \cdot(\hbar . \mathcal{H}) \cdot t \cdot \operatorname{sech}^{4}(x) \tanh (x)-2^{4} \cdot(\hbar . \mathcal{H}) \cdot t \cdot \operatorname{sech}^{2}(x) \\
& \tanh ^{3}(x) . \\
& u_{2}=-2^{4} \cdot\left((\hbar \cdot \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t . \operatorname{sech}^{4}(x) \tanh (x)-2^{4}((\hbar . \mathcal{H}) \\
&+\left.(\hbar . \mathcal{H})^{2}\right) \cdot t \cdot \operatorname{sech}^{2}(x) \tanh ^{3}(x)-2^{5} \cdot(\hbar \cdot \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{8}(x)+ \\
& 2^{5} \cdot 3(\hbar . \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \tanh ^{4}(x)+2^{6} \cdot(\hbar \cdot \mathcal{H})^{2} \cdot t^{2} \cdot \\
& \operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x)
\end{aligned}
$$

$u_{3 .}=(1+(\hbar . \mathcal{H})) \cdot\left[-2^{4} \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t \cdot \operatorname{sech}^{4}(x) \tanh (x)\right.$ $-2^{4} \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) . t . \operatorname{sech}^{2}(x) \cdot \tanh ^{3}(x)-2^{5}(\hbar \mathcal{H})^{2}$. $t^{2} \cdot \operatorname{sech}^{8}(x)+2^{5} \cdot 3 \cdot(\hbar \cdot \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \tanh ^{4}(x)+2^{6}$. $\left.(\hbar . \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x)\right]+(\hbar . \mathcal{H}) \cdot\left[-2^{5} \cdot((\hbar . \mathcal{H})+\right.$ $\left.(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{8}(x)+2^{5} \cdot 3 \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t^{2}$. $\operatorname{sech}^{4}(x) \tanh ^{4}(x)+2^{6} \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{2}(x)$ $\tanh ^{6}(x)+2^{10} \cdot(\hbar \cdot \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{10}(x) \tanh (x)+2^{9} \cdot 5$ $(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x) \cdot \tanh ^{3}(x)-2^{9} \cdot(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x)$ $\tanh ^{7}(x)+2^{9} \cdot 3 \cdot(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{6}(x) \tanh ^{5}(x)-2^{9}$. $\left.(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \tanh ^{9}(x)\right]$.
$u_{4}=(1+\hbar . \mathcal{H}) \cdot u_{3}+\left(\hbar . \mathcal{H}+(\hbar . \mathcal{H})^{2}\right) \cdot\left[2^{6} \cdot\left(\hbar . \mathcal{H}^{2}+(\hbar . \mathcal{H})\right)\right.$. $t^{2} . \operatorname{sech}^{2}(x) \tanh ^{6}(x)-2^{5} .\left(\hbar . \mathcal{H}+(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{8}(x)$. $+2^{5} \cdot 3 .\left(\hbar . \mathcal{H}+(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \tanh ^{4}(x)-2^{9} \cdot(\hbar . \mathcal{H})^{2}$.
$\frac{t^{3}}{3} \cdot \operatorname{sech}^{10}(x) \cdot \tanh (x)+2^{9} \cdot 5 \cdot(\hbar \cdot \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x) \tanh ^{7}(x)+$ $2^{11} .3 .(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \operatorname{sech}^{6}(x) \tanh ^{5}(x)+2^{9} \cdot 5 .(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x)$ $\left.\tanh ^{3}(x)-2^{9}(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \operatorname{sech}^{2}(x) \tanh ^{9}(x)\right]-2^{9} \cdot\left((\hbar . \mathcal{H})^{3}+\right.$ $\left.(\hbar . \mathcal{H})^{4}\right) \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{9}(x)+2^{7} \cdot 37 .\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right)$ $\frac{t^{3}}{3} \cdot \operatorname{sech}^{10}(x) \tanh (x)+2^{9} \cdot 5 .\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x)$. $\tanh ^{3}(x)-2^{9} \cdot 7 \cdot\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x) \tanh ^{7}(x)-2^{10}$
3. $\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{6}(x) \tanh ^{5}(x)-2^{12} \cdot 3 \cdot 5 .(\hbar . \mathcal{H})^{4}$. $\frac{t^{4}}{12} \cdot \operatorname{sech}^{10}(x) \tanh ^{4}(x)-2^{11} \cdot 3 \cdot(\hbar \cdot \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{12}(x) \cdot \tanh ^{2}(x)$ $+2^{12} \cdot(\hbar \cdot \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{14}(x)-2^{12} \cdot 3.5 .(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \operatorname{sech}^{6}(x)$. $\tanh ^{8}(x)-2^{12} \cdot 5^{2} .(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{8}(x) \tanh ^{6}(x)-2^{11} .3$ $(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{4}(x) \tanh ^{10}(x)+2^{12} \cdot(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{2}(x)$. $\tanh ^{12}(x)$.

Therefore the solution series is
$u(x, t)=u_{0}+u_{1}+u_{2}+u_{3}+\cdots \cdots \cdots$
$=2 \operatorname{sech}^{2}(x)-2^{4} .(\hbar . \mathcal{H}) . t \cdot \operatorname{sech}^{4}(x) \tanh (x)-2^{4} .(\hbar . \mathcal{H}) . t$.
$\operatorname{sech}^{2}(x) \cdot \tanh ^{3}(x)-2^{4} \cdot\left((\hbar \cdot \mathcal{H})+(\hbar \cdot \mathcal{H})^{2}\right) \cdot t \cdot \operatorname{sech}^{4}(x) \tanh (x)$ $-2^{4}\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) . t . \operatorname{sech}^{2}(x) \tanh ^{3}(x)-2^{5} .(\hbar . \mathcal{H})^{2} \cdot t^{2}$.
$\operatorname{sech}^{8}(x)+2^{5} \cdot 3(\hbar \cdot \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \tanh ^{4}(x)+2^{6}(\hbar \cdot \mathcal{H})^{2} \cdot t^{2}$. $\operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x)+(1+(\hbar . \mathcal{H})) \cdot\left[-2^{4} \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t\right.$. $\operatorname{sech}^{4}(x) \tanh (x)-2^{4} \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{3}(x)$ $-2^{5} .(\hbar . \mathcal{H})^{2} \cdot t^{2} . . \operatorname{sech}^{8}(x)+2^{5} .3 .(\hbar . \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \tanh ^{4}(x)$ $\left.+2^{6} \cdot(\hbar \cdot \mathcal{H})^{2} \cdot t^{2} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x)\right]+(\hbar \cdot \mathcal{H}) \cdot\left[-2^{5} \cdot((\hbar . \mathcal{H})+\right.$ $\left.(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{8}(x)+2^{5} \cdot 3 \cdot\left((\hbar . \mathcal{H})+(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{4}(x)$. $\tanh ^{4}(x)+2^{6} \cdot\left((\hbar \cdot \mathcal{H})+(\hbar \cdot \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x)+2^{10}$. $(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{10}(x) \tanh (x)+2^{9} \cdot 5 .(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x)$. $\tanh ^{3}(x)-2^{9} .(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x) \tanh ^{7}(x)+2^{9} .3 .(\hbar . \mathcal{H})^{2}$. $\left.\frac{t^{3}}{3} \cdot \operatorname{sech}^{6}(x) \tanh ^{5}(x)-2^{9} \cdot(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \tanh ^{9}(x)\right]+$ $(1+(\hbar . \mathcal{H})) \cdot u_{3}+\left(\hbar . \mathcal{H}+(\hbar . \mathcal{H})^{2}\right) \cdot\left[2^{6} \cdot\left(\hbar . \mathcal{H}+(\hbar . \mathcal{H})^{2}\right) \cdot \mathrm{t}^{2}\right.$. $\operatorname{sech}^{2}(x) \tanh ^{6}(x)-2^{5} \cdot\left(\hbar . \mathcal{H}+(\hbar . \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{8}(x)+2^{5} \cdot 3$
$\left(\hbar \cdot \mathcal{H}+(\hbar \cdot \mathcal{H})^{2}\right) \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \tanh ^{4}(x)-2^{9} \cdot(\hbar \cdot \mathcal{H})^{2} \cdot \frac{\mathrm{t}^{3}}{3}$. $\operatorname{sech}^{10}(x) \cdot \tanh (x)+2^{9} \cdot 5 \cdot(\hbar \cdot \mathcal{H})^{2} \cdot \frac{\mathrm{t}^{3}}{3} \cdot \operatorname{sech}^{4}(\mathrm{x}) \tanh ^{7}(\mathrm{x})+2^{11} \cdot 3$. $(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{6}(x) \tanh ^{5}(x)+2^{9} \cdot 5 \cdot(\hbar \cdot \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x)$ $\left.\tanh ^{3}(x)-2^{9} \cdot(\hbar . \mathcal{H})^{2} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \tanh ^{9}(x)\right]-2^{9} \cdot\left((\hbar . \mathcal{H})^{3}+\right.$ $\left.(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{9}(x)+2^{7} \cdot 37 \cdot\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right)$. $\frac{t^{3}}{3} \cdot \operatorname{sech}^{10}(x) \tanh (x)+2^{9} \cdot 5 \cdot\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x)$. $\tanh ^{3}(x)-2^{9} \cdot 7\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x) \tanh ^{7}(x)-2^{10}$. 3. $\left((\hbar . \mathcal{H})^{3}+(\hbar . \mathcal{H})^{4}\right) \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{6}(x) \tanh ^{5}(x)-2^{12} \cdot 3 \cdot 5 \cdot(\hbar . \mathcal{H})^{4}$. $\frac{t^{4}}{12} \cdot \operatorname{sech}^{10}(x) \tanh ^{4}(x)-2^{11} \cdot 3 .(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{12}(x) \cdot \tanh ^{2}(x)$ $+2^{12} \cdot(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{14}(x)-2^{12} \cdot 3.5 .(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \operatorname{sech}^{6}(x)$. $\tanh ^{8}(x)-2^{12} \cdot 5^{2} \cdot(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{8}(x) \tanh ^{6}(x)-2^{11} .3$. $(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{4}(x) \tanh ^{10}(x)+2^{12} \cdot(\hbar . \mathcal{H})^{4} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{2}(x)$ $\tanh ^{12}(x)+\cdots \cdots \cdots$,

### 3.2 Numerical Scheme by the HPM:

For solving equation (a) by the HPM, we start by making a homotopy
$w: \Omega \times[0,1] \rightarrow \mathbb{R}^{2}$, which satisfies the homotopy equation
$H(w, p)=\mathcal{L}(w)-\mathcal{L}\left(u_{0}\right)+p \mathcal{L}\left(u_{0}\right)+p[\mathcal{N}(w)-f(x, t)]=0$,
where $\quad \mathcal{L}=\frac{\partial}{\partial \mathrm{t}}, \mathcal{N}(w)=6 w w_{\mathrm{x}}+w_{\mathrm{xxx}}, f(x, t)=0 \& p \in[0,1]$.
Then we have, $w_{t}-\left(u_{0}\right)_{t}+p\left(u_{0}\right)_{t}+p\left[6 w w_{x}+w_{x x x}\right]=0$.

Substituting the initial condition in equation (26), we have,

$$
\begin{align*}
& w_{\mathrm{t}}-\left(2 \operatorname{sech}^{2}(x)\right)_{\mathrm{t}}+\mathrm{p}\left(2 \operatorname{sech}^{2}(x)\right)_{\mathrm{t}}+\mathrm{p}\left[6 w w_{\mathrm{x}}+w_{\mathrm{xxx}}\right]=0 \\
& \text { i.e., } w_{\mathrm{t}}-0+\mathrm{p} .0+\mathrm{p}\left[6 w w_{\mathrm{x}}+w_{\mathrm{xxx}}\right]=0 \\
& \text { i.e., } w_{\mathrm{t}}+\mathrm{p}\left[6 w w_{\mathrm{x}}+w_{\mathrm{xxx}}\right]=0 . \quad \ldots \ldots \ldots(27) \tag{27}
\end{align*}
$$

In equation (27), substituting $w=w_{0}+p w_{1}+p^{2} w_{2}+$
$p^{3} w_{3}+\cdots$, we have
$\left(w_{0}+p w_{1}+p^{2} w_{2}+p^{3} w_{3}+\cdots\right)+\mathrm{p}\left[6 .\left(w_{0}+p w_{1}+p^{2} w_{2}+\right.\right.$ $\left.p^{3} w_{3}+\cdots\right) \cdot\left(w_{0}+p w_{1}+p^{2} w_{2}+p^{3} w_{3}+\cdots\right)_{\mathrm{x}}+\left(w_{0}+p w_{1}+\right.$ $\left.p^{2} w_{2}+p^{3} w_{3}+\cdots\right)_{\mathrm{xxx}}$

For simplification we consider $u(x, 0)=w(x, 0)=2 \operatorname{sech}^{2}(x)$

$$
\text { i.e., }\left(w_{0}+p w_{1}+p^{2} w_{2}+p^{3} w_{3}+\cdots \cdots \cdots\right)(x, 0)=2 \operatorname{sech}^{2}(x) .
$$

Which implies that

$$
w_{0}(x, 0)=2 \operatorname{sech}^{2}(x) ; w_{1}(x, 0)=w_{2}(x, 0)=w_{3}(x, 0)=\cdots=0 .
$$

Now, equation (28) can be written as

$$
\begin{aligned}
& p^{0}\left(w_{0}\right)_{t}+p^{1}\left[\left(w_{1}\right)_{t}+6 w_{0}\left(w_{0}\right)_{x}+\left(w_{0}\right)_{x x x}\right]+p^{2}\left[\left(w_{2}\right)_{t}+6 w_{0} .\right. \\
& \left.\left(w_{1}\right)_{x}+6 w_{1}\left(w_{0}\right)_{x}+\left(w_{1}\right)_{x x x}\right]+p^{3}\left[\left(w_{3}\right)_{t}+6 w_{2}\left(w_{0}\right)_{x}+6 w_{1} .\right. \\
& \left.\left(w_{1}\right)_{x}+6 w_{0}\left(w_{2}\right)_{x}+\left(w_{2}\right)_{x x x}\right]+\cdots \cdots \cdots+p^{n}\left[\left(w_{n}\right)_{t}+6 .\right. \\
& \sum_{i=0}^{n-1} \mathrm{w}_{\mathrm{i}}\left(\mathrm{w}_{\mathrm{n}-1-\mathrm{i}}\right)_{\mathrm{x}}+\left(w_{n-1}\right)_{x x x}+\cdots \cdots \cdots=0 .
\end{aligned}
$$

This equation can be represented as:

$$
\begin{gathered}
p^{0}:\left(w_{0}\right)_{t}=0 ; \quad w_{0}(x, 0)=2 \operatorname{sech}^{2}(x) \\
p^{1}:\left(w_{1}\right)_{t}+6 w_{0}\left(w_{0}\right)_{x}+\left(w_{0}\right)_{x x x}=0 ; w_{1}(x, 0)=0 \\
p^{2}:\left(w_{2}\right)_{t}+6 w_{0}\left(w_{1}\right)_{x}+6 w_{1}\left(w_{0}\right)_{x}+\left(w_{1}\right)_{x x x}=0 ; w_{2}(x, 0)=0, \\
p^{3}:\left(w_{3}\right)_{t}+6 w_{0}\left(w_{2}\right)_{x}+6 w_{1}\left(w_{1}\right)_{x}+6 w_{2}\left(w_{0}\right)_{x}+\left(w_{2}\right)_{x x x}=0 ; \\
w_{3}(x, 0)=0,
\end{gathered}
$$

:
$p^{n}:\left(w_{n}\right)_{t}+6 . \sum_{i=0}^{n-1} \mathrm{w}_{\mathrm{i}}\left(\mathrm{w}_{\mathrm{n}-1-\mathrm{i}}\right)_{\mathrm{x}}+\left(w_{n-1}\right)_{x x x}=0 ; w_{n}(x, 0)=0$,

Solving the above equation we can find $w_{0}, w_{1}, w_{2}, w_{3}, \cdots$, as follows:

$$
\begin{aligned}
& w_{0}=2 \operatorname{sech}^{2}(x) ; \\
& w_{1}=2^{4} \cdot t \cdot \operatorname{sech}^{4}(x) \cdot \tanh (x)+2^{4} \cdot t \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{3}(x) ; \\
& w_{2}=2^{5} \cdot 3 \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \cdot \tanh ^{4}(x)-2^{5} \cdot t^{2} \cdot \operatorname{sech}^{8}(x)+2^{6} \cdot t^{2} . \\
& \operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x) ; \\
& w_{3}=-2^{9} \cdot 5 \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x) \cdot \tanh ^{3}(x)-2^{10} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{10}(x) \text {. } \\
& \tanh (x)+2^{9} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x) \cdot \tanh ^{7}(x)-2^{9} \cdot 3 \cdot \frac{t^{3}}{3} . \\
& \operatorname{sech}^{6}(x) \cdot \tanh ^{5}(x)+2^{9} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{9}(x) \text {; } \\
& w_{4}=-2^{12} \cdot 3 \cdot 5 \cdot \frac{\mathrm{t}^{4}}{12} \cdot \operatorname{sech}^{10}(\mathrm{x}) \tanh ^{4}(\mathrm{x})-2^{11} \cdot 3 \cdot \frac{\mathrm{t}^{4}}{12} \cdot \operatorname{sech}^{12}(\mathrm{x}) . \\
& \tanh ^{2}(x)+2^{12} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{14}(x)-2^{12} \cdot 3 \cdot 5 \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{6}(x) . \\
& \tanh ^{8}(x)-2^{12} \cdot 5^{2} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{8}(x) \tanh ^{6}(x)-2^{11} \cdot 3 \cdot \frac{t^{4}}{12} . \\
& \operatorname{sech}^{4}(x) \tanh ^{10}(x)+2^{12} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{2}(x) \tanh ^{12}(x) ; \\
& \text { As a result, the solution series becomes } \\
& u(x, t)=\lim _{\mathrm{p} \rightarrow 1} w(x, t, p)=w_{0}+w_{1}+w_{2}+w_{3}+\cdots \cdots \cdots \\
& =2 \operatorname{sech}^{2}(x)+2^{4} t \operatorname{sech}^{4}(x) \tanh (x)+2^{4} . t \operatorname{sech}^{2}(x) \tanh ^{3}(x) \\
& +2^{5} \cdot 3 \cdot t^{2} \cdot \operatorname{sech}^{4}(x) \cdot \tanh ^{4}(x)-2^{5} \cdot t^{2} \cdot \operatorname{sech}^{8}(x)+2^{6} \cdot t^{2} .
\end{aligned}
$$

$\operatorname{sech}^{2}(x) \cdot \tanh ^{6}(x)-2^{9} \cdot 5 \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{8}(x) \cdot \tanh ^{3}(x)-2^{10} \cdot \frac{t^{3}}{3}$. $\operatorname{sech}^{10}(x) \cdot \tanh (x)+2^{9} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{4}(x) \cdot \tanh ^{7}(x)-2^{9} \cdot 3 \cdot \frac{t^{3}}{3}$.
$\operatorname{sech}^{6}(x) \cdot \tanh ^{5}(x)+2^{9} \cdot \frac{t^{3}}{3} \cdot \operatorname{sech}^{2}(x) \cdot \tanh ^{9}(x)-2^{12} \cdot 3.5$.
$\frac{t^{4}}{12} \cdot \operatorname{sech}^{10}(x) \tanh ^{4}(x)-2^{11} \cdot 3 \frac{t^{4}}{12} \operatorname{sech}^{12}(x) \cdot \tanh ^{2}(x)+2^{12} \cdot \frac{t^{4}}{12}$ $\operatorname{sech}^{14}(x)-2^{12} \cdot 3 \cdot 5 \frac{t^{4}}{12} \operatorname{sech}^{6}(x) \tanh ^{8}(x)-2^{12} \cdot 5^{2} \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{8}(x)$ $\cdot \tanh ^{6}(x)-2^{11} \cdot 3 \cdot \frac{t^{4}}{12} \cdot \operatorname{sech}^{4}(x) \tanh ^{10}(x)+2^{12} \cdot \frac{t^{4}}{12} \operatorname{sech}^{2}(x)$. $\tanh ^{12}(x)+\cdots \ldots \ldots$

## 4 Results And discussion:

We have computed the results of the K-dV equation by both HAM and HPM using our MATLAB routine and have compared the results with the exact solution in the interval $[-5,5]$. In this case we have considered $c=4$ in the exact solution

$$
\begin{equation*}
u(x, t)=\frac{c}{2} \operatorname{sech}^{2}\left(\frac{\sqrt{c}}{2}(x-c t)\right) \tag{b}
\end{equation*}
$$

It is to be noted here that we have considered three cases for both HAM and HPM. Case-I is with three terms $u_{0}, u_{1}, u_{2}$ for HAM, and $w_{0}, w_{1}, w_{2}$ for HPM. Case-II is with four terms $u_{0}, u_{1}, u_{2}, u_{3}$ for HAM, and $w_{0}, w_{1}, w_{2}, w_{3}$ for HPM. And Case-III is with five terms $u_{0}, u_{1}, u_{2}, u_{3}, u_{4}$ for HAM, and $w_{0}, w_{1}, w_{2}, w_{3}, w_{4}$ for HPM.

### 4.1 DISCUSSION OF THE NUMERICAL RESULTS:

Fig. 1 shows that for the product of auxiliary parameter $(\hbar)$ and auxiliary function $(\mathcal{H}(\mathrm{t}))$ equal to-1, i.e., $\hbar \mathcal{H}=-1$, the solution of the K-dV equation (a) obtained by HAM coincides with the solution obtained by HPM and both solutions agree with the exact solution (b) for all the three cases.
Figs. $2,3,5$, and 6 show that for $(\mathcal{H}, \hbar)=(-1,0.5),(-1,-0.5)$, $(-1,0.25),(-0.5,-1),(-0.5,-0.5),(-0.5,0.5),(-0.5,1)$,
$(-0.25,-1),(-0.25,1),(0.5,-1),(0.5,-0.5)(0.5,0.5),(0.5,1)$,
$(1,-0.5),(1,0.25)$ and $(1,0.5)$, HAM gives better solution for large values of time $t$ where as HPM gives more accurate results for small values of $t$ and results from both HAM and HPM agree well with the exact solution (b) for all the three cases.
Fig. 4 shows that for $(\mathcal{H}, \hbar)=(-1,-1),(1,1),\left(-2,-\frac{1}{2}\right)$, $\left(\frac{1}{2}, 2\right),\left(3, \frac{1}{3}\right)$, and $\left(-\frac{1}{3},-3\right)$ HPM solution is more accurate than HAM for all values of $t$ for all the three cases. It is to be noted here that for small values of $t$ results from both HAM and HPM agree well with the exact solution (b) for all the three cases.

### 4.2 ERROR ANALYSIS:

Using our MATLAB routine, we have been computed $l_{2}$-errors of the solution of K-dV equation (a) obtained by both HAM and HPM for all the three cases in different values of time $t$ using the exact solution (b) and the error formula, error $=\sqrt{\sum_{i=1}^{n}\left(f_{e}\left(x_{i}\right)-f_{c}\left(x_{i}\right)\right)^{2}}$, where $f_{\mathrm{e}}, f_{\mathrm{c}}$ are exact and computed solutions, respectively. The results represented in Table $(1-6)$. The results show that for smaller values of time $t$ the error decreases. For all the cases HPM has minimum error for very small values of time $t$.



Fig. 1: $1^{\text {st }}, 2^{\text {nd }}$ and $3^{\text {rd }}$ columns represent the exact solution, HAM solution and HPM solution of the K-dV equation for Case-I, CaseII and Case-III, respectively with regard to the values $(\mathcal{H}, \hbar)=(-1,1),(1,-1),\left(-2, \frac{1}{2}\right),\left(-\frac{1}{2}, 2\right),\left(3,-\frac{1}{3}\right)$, and $\left(\frac{1}{3},-3\right)$ in HAM.



Fig. 2: $1^{\text {st }}, 2^{\text {nd }}$ and $3^{\text {rd }}$ columns represent the exact solution, HAM solution and HPM solution of the K-dV equation for Case-I, Case-II and Case-III, respectively with regard to the values $(\mathcal{H}, \hbar)=(-1,0.5),(1,-0.5),(0.5,-1)$, and $(-0.5,1)$ in HAM.




Fig. 3: $1^{\text {st }}, 2^{\text {nd }}$ and $3^{\text {rd }}$ columns represent the exact solution, HAM solution and HPM solution of the K-dV equation for Case-I, Case-II and Case-III, respectively with regard to the values $(\mathcal{H}, \hbar)=(-1,-0.5),(-0.5,-1),(1,0.5)$, and $(0.5,1)$ in HAM.




Fig. 4: $1^{\text {st }}, 2^{\text {nd }}$ and $3^{\text {rd }}$ columns represent the exact solution, HAM solution and HPM solution of the K-dV equation for Case-I, Case-II and Case-III, respectively with regard to the values $(\mathcal{H}, \hbar)=(-1,-1),(1,1),\left(-2,-\frac{1}{2}\right),\left(\frac{1}{2}, 2\right),\left(3, \frac{1}{3}\right)$, and $\left(-\frac{1}{3},-3\right)$ in HAM.




Fig. 5: $1^{\text {st }}, 2^{\text {nd }}$ and $3^{\text {rd }}$ columns represent the exact solution, HAM solution and HPM solution of the K-dV equation for Case-I, Case-II and Case-III, respectively with regard to the values $(\mathcal{H}, \hbar)=(-0.5,0.5),(0.5,-0.5),(-1,0.25)$, and $(-0.25,1)$ in HAM.



Fig. 6: $1^{\text {st }}, 2^{\text {nd }}$ and $3^{\text {rd }}$ columns represent the exact solution, HAM solution and HPM solution of the K-dV equation for Case-I, Case-II and Case-III, respectively with regard to the values $(\mathcal{H}, \hbar)=(-0.5,-0.5),(0.5,0.5),(1,0.25)$, and $(-0.25,-1)$ in HAM.

## TABLE 1

The $l_{2}$-error for HAM and HPM solution for all the three cases for the $\operatorname{values}(\mathcal{H}, \hbar)=(-1,1),(1,-1),\left(-2, \frac{1}{2}\right)$, $\left(-\frac{1}{2}, 2\right),\left(3,-\frac{1}{3}\right)$, and $\left(\frac{1}{3},-3\right)$ considered in HAM.

| Time t | Method | Case-I | Case-II | Case-III |
| :---: | :---: | :---: | :---: | :---: |
| $2^{0}$ | HAM | 62.718199989 | 193.753896419 | 577.584811761 |
|  | HPM | 62.718199989 | 193.753896419 | 577.584811761 |
| $2^{-1}$ | HAM | 14.841161323 | 22.282477407 | 32.612810041 |
|  | HPM | 14.841161323 | 22.282477407 | 32.612810041 |
| $2^{-2}$ | HAM | 2.631481915 | 1.973099259 | 1.438693389 |
|  | HPM | 2.631481915 | 1.973099259 | 1.438693389 |
| $2^{-3}$ | HAM | 0.369167277 | 0.138889796 | 0.050721759 |
|  | HPM | 0.369167277 | 0.138889796 | 0.050721759 |
| $2^{-4}$ | HAM | 0.047605331 | 0.008966728 | 0.001638531 |
|  | HPM | 0.047605331 | 0.008966728 | 0.001638531 |
| $2^{-5}$ | HAM | 0.005998090 | 0.000565086 | 0.000005164 |
|  | HPM | 0.005998090 | 0.000565086 | 0.000005164 |
| $2^{-6}$ | HAM | 0.000751258 | 0.000035392 | 0.000000005 |
|  | HPM | 0.000751258 | 0.000035392 | 0.000000005 |

TABLE 2
The $l_{2}$-error for HAM and HPM solution for all the three cases for the values $(\mathcal{H}, \hbar)=(-1,0.5),(1,-0.5)$, $(0.5,-1)$, and $(-0.5,1)$ considered in HAM.

| Time t | Method | Case-I | Case-II | Case-III |
| :---: | :---: | :---: | :---: | :---: |
| $2^{0}$ | HAM | 19.862446963 | 33.769828542 | 70.544316243 |
|  | HPM | 62.718199989 | 193.753896419 | 577.584811761 |
| $2^{-1}$ | HAM | 7.061387460 | 7.276843876 | 8.137547253 |
|  | HPM | 14.841161323 | 22.282477407 | 32.612810041 |
| $2^{-2}$ | HAM | 2.200176138 | 1.621499716 | 1.058254270 |
|  | HPM | 2.631481915 | 1.973099259 | 1.438693389 |
| $2^{-3}$ | HAM | 0.744438007 | 0.441997160 | 0.253230788 |
|  | HPM | 0.369167277 | 0.138889796 | 0.050721759 |
| $2^{-4}$ | HAM | 0.311492307 | 0.165048927 | 0.089137213 |
|  | HPM | 0.047605331 | 0.008966728 | 0.001638531 |
| $2^{-5}$ | HAM | 0.147264288 | 0.074827306 | 0.038415451 |
|  | HPM | 0.005998090 | 0.000565086 | 0.000005164 |
| $2^{-6}$ | HAM | 0.072537104 | 0.036419152 | 0.018341876 |
|  | HPM | 0.000751258 | 0.000035392 | 0.000000005 |
| $2^{-7}$ | HAM | 0.036130509 | 0.018084119 | 0.009058835 |
|  | HPM | 0.000093954 | 0.000002213 | 0.000000000 |

## TABLE 3

The $l_{2}$-error for HAM and HPM solution for all the three cases for the values $(\mathcal{H}, \hbar)=(-1,-0.5),(1,0.5),(0.5,1)$, and $(-0.5,-1)$ considered in HAM.

| Time t | Method | Case-I | Case-II | Case-III |
| :---: | :---: | :---: | :---: | :---: |
| $2^{2}$ | HAM | 27.284422088 | 67.810320782 | 196.506643920 |
|  | HPM | 62.718199989 | 193.753896419 | 577.584811761 |
| $2^{-1}$ | HAM | 14.163199159 | 25.303782693 | 46.794994445 |
|  | HPM | 14.841161323 | 22.282577407 | 32.612810041 |
| $2^{-2}$ | HAM | 8.954242418 | 13.791224761 | 21.410765669 |
|  | HPM | 2.631481915 | 1.973099259 | 1.438693389 |
| $2^{-3}$ | HAM | 4.984952844 | 7.508981031 | 11.326092940 |
|  | HPM | 0.369167277 | 0.138889796 | 0.050721759 |
| $2^{-4}$ | HAM | 2.570545445 | 3.859229076 | 5.795564534 |
|  | HPM | 0.047605331 | 0.008966728 | 0.001638531 |
| $2^{-5}$ | HAM | 1.295560069 | 1.943748812 | 2.916426582 |
|  | HPM | 0.005998090 | 0.000565086 | 0.000005164 |
| $2^{-6}$ | HAM | 0.649083119 | 0.973675208 | 1.460612067 |
|  | HPM | 0.000751258 | 0.000035392 | 0.000000162 |
| $2^{-7}$ | HAM | 0.324704987 | 0.487063779 | 0.730608039 |
|  | HPM | 0.000093954 | 0.000002213 | 0.000000005 |
| $2^{-8}$ | HAM | 0.162372939 | 0.243560195 | 0.365341838 |
|  | HPM | 0.000011746 | 0.000000138 | 0.000000000 |
| $2^{-9}$ | HAM | 0.081189026 | 0.121783637 | 0.1826756483 |
|  | HPM | 0.000001468 | 0.000000009 | 0.0000000000 |

## TABLE 4

The $l_{2}$-error for HAM and HPM solution for all the three cases for the values $(\mathcal{H}, \hbar)=(-1,-1),(1,1),\left(-2,-\frac{1}{2}\right),\left(\frac{1}{2}, 2\right)$, $\left(3, \frac{1}{3}\right)$, and $\left(-\frac{1}{3},-3\right)$ considered in HAM.

| Time t | Method | Case-I | Case-II | Case-III |
| :---: | :---: | :---: | :---: | :---: |
| $2^{0}$ | HAM | 81.739945215 | 339.130355542 | 1839.752583995 |
|  | HPM | 62.718199989 | 193.753896419 | 577.584811761 |
| $2^{-1}$ | HAM | 32.151611344 | 91.105747660 | 288.842157676 |
|  | HPM | 14.841161323 | 22.282477407 | 32.612810041 |
| $2^{-2}$ | HAM | 16.899705653 | 37.326398737 | 83.363702134 |
|  | HPM | 2.631481915 | 1.973099259 | 1.438693389 |
| $2^{-3}$ | HAM | 8.984843690 | 18.393021229 | 37.582020926 |
|  | HPM | 0.369167277 | 0.1388897959 | 0.050721759 |
| $2^{-4}$ | HAM | 4.585151957 | 9.222230737 | 18.532962159 |
|  | HPM | 0.047605331 | 0.008966728 | 0.001638531 |
| $22^{-5}$ | HAM | 2.305127771 | 4.616711291 | 9.244127321 |
|  | HPM | 0.005998090 | 0.000565086 | 0.000005164 |
| $2^{-6}$ | HAM | 1.154164220 | 2.309134279 | 4.619595362 |
|  | HPM | 0.000751258 | 0.000035392 | 0.000000162 |
| $2^{-7}$ | HAM | 0.577283143 | 1.154666980 | 2.309499457 |
|  | HPM | 0.000093954 | 0.000002213 | 0.000000005 |
| $2^{-8}$ | HAM | 0.288666731 | 0.577346049 | 1.154712774 |
|  | HPM | 0.000011757 | 0.000000138 | 0.000000000 |

TABLE 5
The $l_{2}$-error for HAM and HPM solution for all the three cases for the values $(\mathcal{H}, \hbar)=(-0.5,0.5),(0.5,-0.5)$, $(-1,0.25)$, and $(-0.25,1)$ considered in HAM.

| Time $t$ | Method | Case-I | Case-II | Case-III |
| :---: | :---: | :---: | :---: | :---: |
| $2^{0}$ | HAM | 10.354100581 | 12.811010909 | 16.311437360 |
|  | HPM | 62.718199990 | 193.753896420 | 577.584811761 |
| $2^{-1}$ | HAM | 5.996894130 | 5.895856476 | 5.614762687 |
|  | HPM | 14.841161323 | 22.282477407 | 32.612810041 |
| $2^{-2}$ | HAM | 2.857547504 | 2.412555823 | 2.032031850 |
|  | HPM | 2.631481915 | 1.973099259 | 1.438693389 |
| $2^{-3}$ | HAM | 1.341834764 | 1.045717677 | 0.819790923 |
|  | HPM | 0.369167277 | 0.138889796 | 0.050721759 |
| $2^{-4}$ | HAM | 0.655278965 | 0.496609591 | 0.377323721 |
|  | HPM | 0.047605331 | 0.008966728 | 0.001638531 |
| $2^{-5}$ | HAM | 0.325493193 | 0.244771618 | 0.184203194 |
|  | HPM | 0.005998090 | 0.000565086 | 0.000005164 |
| $22^{-6}$ | HAM | 0.16247190 | 0.121935644 | 0.091530325 |
|  | HPM | 0.000751258 | 0.000035392 | 0.000000162 |
| $22^{-7}$ | HAM | 0.081201409 | 0.060911280 | 0.045693301 |
|  | HPM | 0.000093954 | 0.000002213 | 0.000000005 |

## TABLE 6

The $l_{2}$-error for HAM and HPM solution for all the three cases for the values $(\mathcal{H}, \hbar)=(-0.5,-0.5),(0.5,0.5),(1,0.25)$, and $(-0.25,-1)$ considered in HAM.

| Time t | Method | Case-I | Case-II | Case-III |
| :---: | :---: | :---: | :---: | :---: |
| $2^{0}$ | HAM | 12.416474038 | 20.157536021 | 33.814347982 |
|  | HPM | 62.718199989 | 193.753896419 | 577.584811761 |
| $2^{-1}$ | HAM | 8.906696778 | 11.425235031 | 15.069595136 |
|  | HPM | 14.841161323 | 22.282477407 | 32.612810041 |
| $2^{-2}$ | HAM | 6.150968963 | 7.627559517 | 9.519356544 |
|  | HPM | 2.631481915 | 1.973099259 | 1.438693389 |
| $2^{-3}$ | HAM | 3.456631905 | 4.307124161 | 5.373854083 |
|  | HPM | 0.369167277 | 0.138889796 | 0.050721759 |
| $2^{-4}$ | HAM | 1.784580840 | 2.228791702 | 2.784413109 |
|  | HPM | 0.047605331 | 0.008966728 | 0.001638531 |
| $2^{-5}$ | HAM | 0.899633476 | 1.124292643 | 1.405158129 |
|  | HPM | 0.005998090 | 0.000565086 | 0.000005164 |
| $2^{-6}$ | HAM | 0.450744665 | 0.563399447 | 0.704223014 |
|  | HPM | 0.000751258 | 0.000035392 | 0.000000162 |
| $2^{-7}$ | HAM | 0.225488642 | 0.281856870 | 0.352317790 |
|  | HPM | 0.000093954 | 0.000002213 | 0.000000005 |

## 6 CONCLUSION:

In this study, both HAM and HPM have been applied to solve the K-dV equation (a) and compared all the results obtained by these two methods with exact solution (b). From Fig. 1 and Table 1 we see that the HAM and HPM solution coincide when the values of $\mathcal{H}$ and $\hbar$ are taken in such a way that $\mathcal{H} \hbar=-1$. For the other setting the HPM gives better results than HAM because the initial approximation is chosen properly. We hope the study will be helpful for further studies of HAM and HPM for other differential equations.

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## Ethical statement:

We have not intentionally engaged in or participate in any form of malicious harm to another person or animal. We have no conflict of interest.

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